

Matrix factorisations and Quantum Error Correcting Codes

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An important conceptual insight from modern logic is that there is a dynamical process "lying behind" function composition. This process is called computation. Taking our cue from the logicians we should, as mathematicians, attempt to find some subatomic structure inside the "atomic" operation $(g, f) \mapsto g \circ f$. Today I'll explain one approach to doing just that in the context of hypersurface singularities.

1. Matrix factorisations

Let k be a commutative ring. A polynomial $W \in k[\underline{x}] = k[x_1, \dots, x_n]$ is a potential if

(i) $\partial_{x_1} W, \dots, \partial_{x_n} W$ is a quasi-regular sequence

(ii) $k[\underline{x}]/(\partial_{x_1} W, \dots, \partial_{x_n} W)$ is a f.g. free k -module

(iii) $H^i(k(\partial_{x_1} W, \dots, \partial_{x_n} W)) = 0 \quad i \neq 0 \quad (\text{note (iii)} \Rightarrow \text{(i)})$

Example $W \in \mathbb{C}[x_1, \dots, x_n]$ with $\dim(k[\underline{x}]/(\partial_{x_1} W, \dots, \partial_{x_n} W)) < \infty$,
e.g. any isolated hypersurface singularity.

Remark If $(k[\underline{x}], W), (k[\underline{y}], V)$ are potentials so are $(k[\underline{x}], -W), (k[\underline{x}, \underline{y}], W + V)$.

Defⁿ (Eisenbud) A matrix factorisation (MF) of $W \in k[\underline{x}]$ is a \mathbb{Z}_2 -graded free $k[\underline{x}]$ -module (possibly infinite rank) $X = X^0 \oplus X^1$ and a $k[\underline{x}]$ -linear odd map $d_X : X \rightarrow X$ such that $d_X^2 = W \cdot 1_X$

$$d_X = \begin{pmatrix} X^0 & X^1 \\ 0 & A \\ B & 0 \end{pmatrix} \quad AB = BA = W \cdot I$$

A morphism $f : (X, d_X) \rightarrow (Y, d_Y)$ is a $k[\underline{x}]$ -linear even map s.t. $d_Y f = f d_X$.

(2)

There are some triangulated categories associated to a potential

$$\mathcal{T} = \text{HMF}(k[x], W)$$

Homotopy category of all MFs

$$\mathcal{T}' = \text{hmf}(k[x], W)^\oplus$$

Full subcategory of direct summands of $Y \in \text{hmf}$

$$S = \text{hmf}(k[x], W)$$

Kanobi completion

Homotopy category of f. rank MFs

Defⁿ The bicategory \mathcal{LG}_k has potentials as objects and

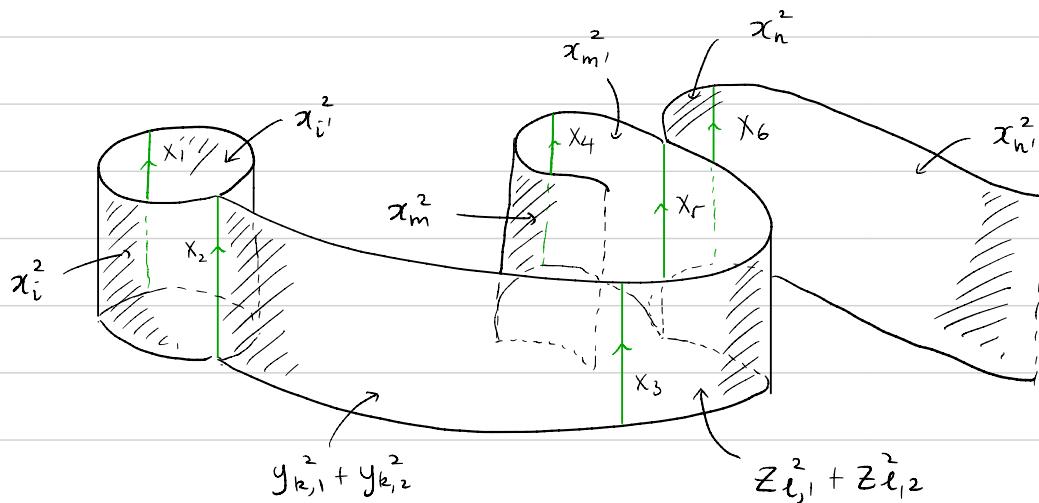
[studied by Lazariu-McNamee,
Khovanov-Rozansky]

$$\mathcal{LG}_k((k[x], W), (k[y], V)) := \text{hmf}(k[x, y], V - W)^\oplus$$

Theorem The bicategory \mathcal{LG}_k

- has adjoints for 1-morphisms (Carqueville-Murfet '16)
- is symmetric monoidal with duals (Carqueville-Montoya '18)
- is a pivotal superbicategory (Godfrey-Murfet '18, '22)

This means \mathcal{LG}_k determines a 2D TQFT of a kind which can evaluate bordisms like :



2. Composition

What has this got to do with "subatomic" structure in composition? Like any bicategory $\mathcal{L}\mathcal{G}_k$ has composition $(Y, X) \mapsto Y \circ X$ for 1-morphisms. Given potentials $W(x), V(y), U(z)$

$$\begin{array}{ccc}
 & (Y, X) & \\
 \text{hmf}(k[y, z], U - V) \times \text{hmf}(k[x, y], V - W) & \xrightarrow{\quad ? \quad} & \text{hmf}(k[x, z], U - W) \\
 \downarrow & \dashrightarrow & \downarrow \\
 \text{hmf}(k[y, z], U - V)^{\oplus} \times \text{hmf}(k[x, y], V - W)^{\oplus} & & \text{hmf}(k[x, z], U - W)^{\oplus} \\
 \mathcal{L}\mathcal{G}(V, U) & \downarrow & \mathcal{L}\mathcal{G}(W, V) & \downarrow & \mathcal{L}\mathcal{G}(W, U) \\
 \text{HMF}(k[y, z], U - V) \times \text{HMF}(k[x, y], V - W) & \xrightarrow{\quad} & \text{HMF}(k[x, z], U - W) & \xrightarrow{Y \otimes X} & \\
 (Y, X) & & (Y \otimes_{k[y]} X, d_Y \otimes 1 + 1 \otimes d_X) & &
 \end{array}$$

We see that $(Y, X) \mapsto Y \circ X$ is not so simple, if we want a finite-rank answer. It involves

1. Write down $Y \otimes X$

2. Invert the equivalence

$$\text{hmf}(k[x, z], U - W) \xrightleftharpoons[\psi]{\cong} \text{hmf}(k[x, z], U - W)^{\oplus} \\
 (Z, e) \qquad \qquad \qquad Y \otimes X$$

3. Try to split e in finite rank MFs

$$\begin{array}{ccccc}
 e & \hookrightarrow & Z & \xrightleftharpoons[\cong]{F} & Y \otimes X \\
 f \uparrow g & & \downarrow c & & \\
 Y \circ X & \dashrightarrow & & & \\
 \text{hmf} & & \text{HMF} & & \text{What is this?}
 \end{array}$$

3. Cut

A supercategory is a category \mathcal{T} and functor $F: \mathcal{T} \rightarrow \mathcal{T}$ with $F^2 = \text{id}$ satisfying some axioms.

Any hmf or HMF is a supercategory with $F(X) = X[1]$. For $n \geq 0$ let C_n denote the Clifford algebra generated by $\gamma_1, \dots, \gamma_n, \gamma_1^+, \dots, \gamma_n^+$ of odd degree s.t.

$$\gamma_i \gamma_j + \gamma_j \gamma_i = 0, \quad \gamma_i^+ \gamma_j^+ + \gamma_j^+ \gamma_i^+ = 0, \quad \gamma_i \gamma_j^+ + \gamma_j^+ \gamma_i = \delta_{ij}.$$

Example (i) $S_n := \Lambda(k\mathbb{O}_1 \oplus \dots \oplus k\mathbb{O}_n)$ is a left C_n -module, γ_i acts as $\mathbb{O}_i^* \circ (-)$, γ_i^+ acts as $\mathbb{O}_i \wedge (-)$. Also $C_n \cong \text{End}_k(S_n)$.

(ii) $S_{m,n} := S_m \otimes_k S_n^*$ is a C_m - C_n -bimodule

Defⁿ Let \mathcal{T} be an idempotent complete supercategory. The Clifford thickening \mathcal{T}^\bullet of \mathcal{T} has as objects tuples (X, n, ρ) consisting of $X \in \text{ob}(\mathcal{T})$ and a left C_n -module structure $\rho: C_n \otimes_k X \rightarrow X$. A morphism $(X, n, \rho) \rightarrow (Y, m, \tau)$ is a morphism of C_m -modules $S_{m,n} \otimes_{C_n} X \rightarrow Y$.

Lemma $\mathcal{T} \xrightarrow{\cong} \mathcal{T}^\bullet$, $X \mapsto (X, 0, \text{id})$.

Proof C_n is Morita trivial, so if $(X, n, \rho) \in \mathcal{T}^\bullet$ there exists $\tilde{X} \in \mathcal{T}$ with $X \cong S_n \otimes_k \tilde{X}$ as C_n -modules. Then $(X, n, \rho) \cong (\tilde{X}, 0, \text{id})$ in \mathcal{T}^\bullet as

$$\begin{aligned} (\tilde{X}, 0, \text{id}) &\longrightarrow (X, n, \rho) \\ \text{is } S_n \otimes_k \tilde{X} &\longrightarrow X. \quad \square \end{aligned}$$

An object X of \mathcal{T}^\bullet specifies an object \tilde{X} of \mathcal{T} indirectly via $\{\gamma_i, \gamma_i^+\}_{i=1}^n$. You may need to compute in order to actually extract \tilde{X} .

Def The superbicategory \mathcal{C} has the same objects as \mathcal{LG} and

$$\mathcal{C}((k[x], w), (k[y], v)) := \left(\text{hmf}(k[x, y], V - W)^\omega \right)^*$$

finite rank MF + structure

with an explicit composition

$$\mathcal{C}(v, u) \times \mathcal{C}(w, v) \longrightarrow \mathcal{C}(w, u) \quad I = (\partial_{y_1} v, \dots, \partial_{y_m} v)$$

$$(y, x) \longmapsto (y|x := \frac{y \otimes k[y] X}{I(y \otimes k[y] X)})$$

$\nearrow \{x_i, x_i^+\}_{i=1}^n$

we call this
the cut operation

$$x_i = [dy + dx, \partial_{y_i} v] =: At_i$$

$$x_i^+ = -\partial_{y_i}(dx) - \frac{1}{2} \sum_q \partial_{y_q} \partial_{y_i}(v) At_q.$$

Theorem There are isomorphisms of Clifford modules (M'18)

$$Y|X \cong \Lambda(k0_1 \oplus \dots \oplus k0_m) \otimes_k (Y \otimes_{k[y]} X) \quad \text{hmf}(V-W)$$

$$\overset{\circ}{\cup} x_i, x_i^+ \quad \overset{\circ}{\cup} \partial_i^+, \partial_i \quad \swarrow \quad \downarrow$$

and an equivalence $\mathcal{LG} \cong \mathcal{C}$.

$$\text{hmf}^\omega(V-W) \xrightarrow{\cong} \text{hmf}^\oplus(V-W)$$

$$\cong \downarrow \quad \parallel$$

$$\mathcal{C}(w, v) := \text{hmf}^\omega(V-W) \dashleftarrow \mathcal{LG}(w, v)$$

$$e = x_1 \dots x_n x_n^+ \dots x_1^+$$

Inside the composition is a picture of pumping energy out of a fermionic system:

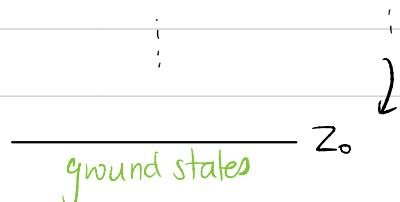
$$\hookrightarrow Y|X \xleftarrow{\quad} Y \otimes X$$

$$x_1 \dots x_{n-1} x_{n-1}^+ \dots x_1^+ \hookrightarrow z_{n-1}$$



$$y_0 X \xleftarrow{\quad} \text{HMF}$$

$$\text{hmf}$$



4. Quantum wider

To ascertain how seriously we should take the analogy to pumping energy out of a physical system, we can consider the simplest possible example: composition of a chain of identity 1-cells on $(\mathbb{C}[x], x^2) \in \mathcal{LG}_C$

$$x_n^2 \xleftarrow{\Delta_n} x_{n-1}^2 \xleftarrow{\quad \cdots \quad} x_2^2 \xleftarrow{\Delta_2} x_1^2$$

$$\Delta_i \in \text{hmf}(\mathbb{C}[x_i, x_{i-1}], x_i^2 - x_{i-1}^2)$$

$$\Delta_i = \Lambda(\mathbb{C}\Psi_i) \otimes_{\mathbb{C}} \mathbb{C}[x] = \mathbb{C}[x]1 \oplus \mathbb{C}[x]\Psi_i$$

$$d_{\Delta_i} = (x_i - x_{i-1})\Psi_i^* + (x_i + x_{i-1})\Psi_i$$

$$= \begin{pmatrix} 1 & \Psi_i \\ 0 & x_i - x_{i-1} \\ x_i + x_{i-1} & 0 \end{pmatrix}$$

We have, since $\mathbb{C}[x]/(\partial_x(x^2)) = \mathbb{C}$

γ, γ^\dagger

$$\begin{aligned} \Delta_n | \Delta_{n-1} | \cdots | \Delta_2 &= \Delta_n \otimes_{\mathbb{C}[x_{n-1}]} \mathbb{C} \otimes_{\mathbb{C}[x_{n-1}]} \Delta_{n-1} \otimes \cdots \\ &\quad \cdots \otimes_{\mathbb{C}[x_2]} \mathbb{C} \otimes_{\mathbb{C}[x_2]} \Delta_2 \\ &\cong \mathbb{C}[x_n, x_1] \otimes_{\mathbb{C}} \left(\Lambda^k \Psi_n \otimes_{\mathbb{C}} \Lambda^k \Psi_{n-1} \otimes \cdots \otimes \Lambda^k \Psi_2 \right) \\ &\cong \mathbb{C}[x_n, x_1] \otimes_{\mathbb{C}} \Lambda(k\Psi_n \oplus \cdots \oplus k\Psi_2) \end{aligned}$$

For each cut we have Clifford operators γ, γ^\dagger . For the Δ_n / Δ_{n-1} cut

$$\begin{aligned} \gamma &= -\frac{1}{2}(\Psi_{n-1}^* + \Psi_{n-1} + \Psi_n - \Psi_n^*) \\ \gamma^\dagger &= -\frac{1}{2}(\Psi_{n-1}^* + \Psi_{n-1} - \Psi_n - \Psi_n^*) \end{aligned} \quad \left\{ \text{notice there are } x\text{-free} \right.$$

(meaning $\Psi_{n-1} \wedge (-)$)
meaning $\Psi_{n-1}^\dagger \perp (-)$

According to our earlier prescription the process of computing $\Delta_n \circ \dots \circ \Delta_2$ is

$$\gamma\gamma^\dagger \underset{\substack{\uparrow \\ \downarrow}}{\mathcal{G}} \Delta_n | \Delta_{n-1} | \dots | \Delta_2 = z_n$$

$$\text{Im}(\gamma\gamma^\dagger) = \text{Ker}(\gamma) = \text{Im}(\gamma^\dagger) = z_{n-1}$$

$$\begin{matrix} \uparrow & & & & (7.1) \\ \downarrow & & & & \\ \vdots & & & & \end{matrix}$$

$$\Delta_n \circ \dots \circ \Delta_2 \quad (= \Delta \text{ again})$$

So what is $\text{Ker}(\gamma)$?

$$-2\gamma = \psi_{n-1}^* + \psi_n + \psi_n - \psi_n^* \underset{\substack{\uparrow \\ \downarrow}}{\mathcal{G}} \Lambda k \psi_n \otimes \Lambda k \psi_{n-1}$$

$$\begin{aligned} \therefore -2\gamma(1 + \psi_n \psi_{n-1}) &= (\psi_{n-1} + \psi_n)(1) + (\psi_{n-1}^* - \psi_n^*)(\psi_n \psi_{n-1}) \\ &= \psi_{n-1} + \psi_n + [-\psi_n - \psi_{n-1}] = 0 \quad (*) \end{aligned}$$

$$-2\gamma(\psi_n + \psi_{n-1}) = 1 + \psi_{n-1}\psi_n + \psi_n\psi_{n-1} - 1 = 0$$

By dimension count $\dim \text{Ker}(\gamma \underset{\substack{\uparrow \\ \downarrow}}{\mathcal{G}} \Lambda k \psi_n \otimes \Lambda k \psi_{n-1}) = 2$ so this is it. But what is going on here?

And how do we continue this to compute $\Delta_n \circ \dots \circ \Delta_2$? Physicists recognise (*) as Bell states or maximally entangled states. They write $\Lambda \mathbb{C}^4 = \mathbb{C}|00\rangle \oplus \mathbb{C}|11\rangle = \mathbb{C}^2$ and

$$1 + \psi_n \psi_{n-1} \text{ as } |00\rangle + |11\rangle$$

$$\psi_n + \psi_{n-1} \text{ as } |10\rangle + |01\rangle$$

Computing $\Delta_n \circ \Delta_{n-1}$ from $\Delta_n | \Delta_{n-1}$ consists in projecting onto these entangled states.

To cut a long story short, the procedure (7.1) rediscovers the error correction process for a particular kind of quantum error correcting code on $\mathcal{H} = \mathbb{C}^2 \otimes \cdots \otimes \mathbb{C}^2$ ($n-1$ copies) called a stabiliser code. With $X|0\rangle = |1\rangle$, $X|1\rangle = |0\rangle$ and $X_i := |0\rangle\langle 0| \otimes \cdots \otimes X \otimes \cdots \otimes |0\rangle\langle 0|$ the code is

$$S = \{X_n X_{n-1}, X_{n-1} X_{n-2}, \dots, X_3 X_2\} \subseteq \text{Aut}_{\mathbb{C}}(\mathcal{H})$$

and the codespace C is the joint $+/-$ -eigenspace of all operators in S . This code is sometimes called a quantum wire (or Majorana chain) and it is closely related to the Ising model.

Lemma The idempotent e computing $\Delta_n \circ \cdots \circ \Delta_1$ from $\Delta_n / \cdots / \Delta_1$ is

$$P = \prod_{i=2}^{n-1} \frac{1 + X_{i+1} X_i}{2} \quad \mathcal{Y} \mathcal{Y}^\dagger = \frac{1 + X_n X_{n-1}}{2}$$

the standard projector for the stabiliser code, hence

$$\Delta_n \circ \cdots \circ \Delta_1 \cong (\mathbb{C}[x_n, x_1]) \otimes_{\mathbb{C}} C \quad \text{codespace}$$

where a \mathbb{C} -basis for C is given by the entangled states

$$|+\cdots+\rangle \pm |- \cdots -\rangle \quad |\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$$

To return to our original theme: the cut operation $(Y, X) \mapsto Y/X$ followed by "eliminating" the Clifford action is our answer to the "subatomic" structure lying inside the atomic operation $(Y, X) \mapsto Y \circ X$. In the simplest case this elimination process is identical with the process of error-correction in a particular well-known stabiliser code, which can be viewed as "pumping out entropy"!

We expect that following this recipe for equivariant MFs will be a new source of error correcting codes.

Theorem Suppose $G \subset k[\underline{y}] = k[y_1, \dots, y_m]$ is a finite group (9)
 acting so V is G -invariant and $g \cdot \partial y_i V \in \text{span}\{\partial y_j V\}_{j=1}^m$.
 Let F_G be the \mathbb{Z}_2 -graded G -rep ($I = (\partial y_1 V, \dots, \partial y_m V)$)

$$F_G \otimes_k k[\underline{y}] / I \cong I / I^2 [1]$$

↑ conormal bundle CrisW

Then for G -equivariant Y, X

$$(Y|X)^G \cong \left[\wedge F_G \otimes_k (Y \otimes_{k[\underline{y}]} X) \right]^G$$

$$\cup$$

$$(Y \otimes_{k[\underline{y}]} X)^G$$